

HOMEWORK #6:

5.10, 6.2, 6.3, 6.4, 6.5

5.10 q, v

- density: $N_i + N_e$
- K temp: $T_i \approx T_e = T$
- Prod $\propto q^2 |dv/dt|^2$

Can't neglect ions, eqn (5.82)

a) $N_D = N_H \Rightarrow N_D = N = N_H$
 $N_e = \frac{3}{2} N_i$

$$P_{cyc} \propto N \left(\frac{q^4 B^2}{m^2} \right) (kT)$$

Deuterium:

$$P_{cyc} \propto \frac{N}{2} \left(\frac{(1) B^2}{(3.3446 \times 10^{-27} \text{ kg})^2} \right) kT$$

Helium

$$P_{cyc} \propto \frac{N}{2} \left(\frac{2^4 B^2}{(5.0084 \times 10^{-27} \text{ kg})^2} \right) kT$$

$$P_{cyc, ions} = \frac{N}{2} B^2 kT \left(\frac{1}{1.187 \times 10^{-53}} \right) + \frac{16}{2.5084 \times 10^{-53}}$$

$$= 3.610 \times 10^{53} B^2 \cdot kT \cdot N$$

electrons

$$P_{cyc} \propto \frac{3}{2} N \left(\frac{B^2}{(9.1095 \times 10^{-31} \text{ kg})^2} \right) kT$$

$$P_{cyc} \propto \frac{3}{2} N B^2 kT (1.20507 \times 10^{60}) = 1.8076 \times 10^{60} N B^2 kT$$

$$\frac{P_e}{P_{ions}} = \left(\frac{N k T B^2}{N k T B^2} \right) \left(\frac{1.8076 \times 10^{60}}{3.610 \times 10^{53}} \right) = 5.0072 \times 10^6$$

$\therefore P_{cyc, ions}$ is 5.0×10^6 times smaller

b) $P_{cyc} = (6.3 \times 10^{-20} \frac{J}{eV \cdot F^2}) (6 \text{ Tesla})^2 (80000 \text{ eV}) (2 \times 10^{20} \frac{1}{m^3}) (3)$

$$P_{cyc} = 1.089 \times 10^8 \frac{W}{m^3}$$

$\sigma N_D q$

$$P_{fu} = \langle \sigma v \rangle N_D N_H Q$$

$$= (1.27 \times 10^{-22}) \frac{m^3}{s} (2 \times 10^{20} \frac{1}{m^3})^2 (18.3 \text{ MeV}) \times \left(\frac{1 \times 10^6 \cdot 1.6022 \times 10^{-19} J}{1 \text{ MeV}} \right)$$

$$P_{fu} = 1.49 \times 10^7 \frac{W}{m^3}$$

P_{cyc} is 7.3 times greater than P_{fu} . This means that radiation reflecting on walls can not be neglected because it is a significant amount in break even

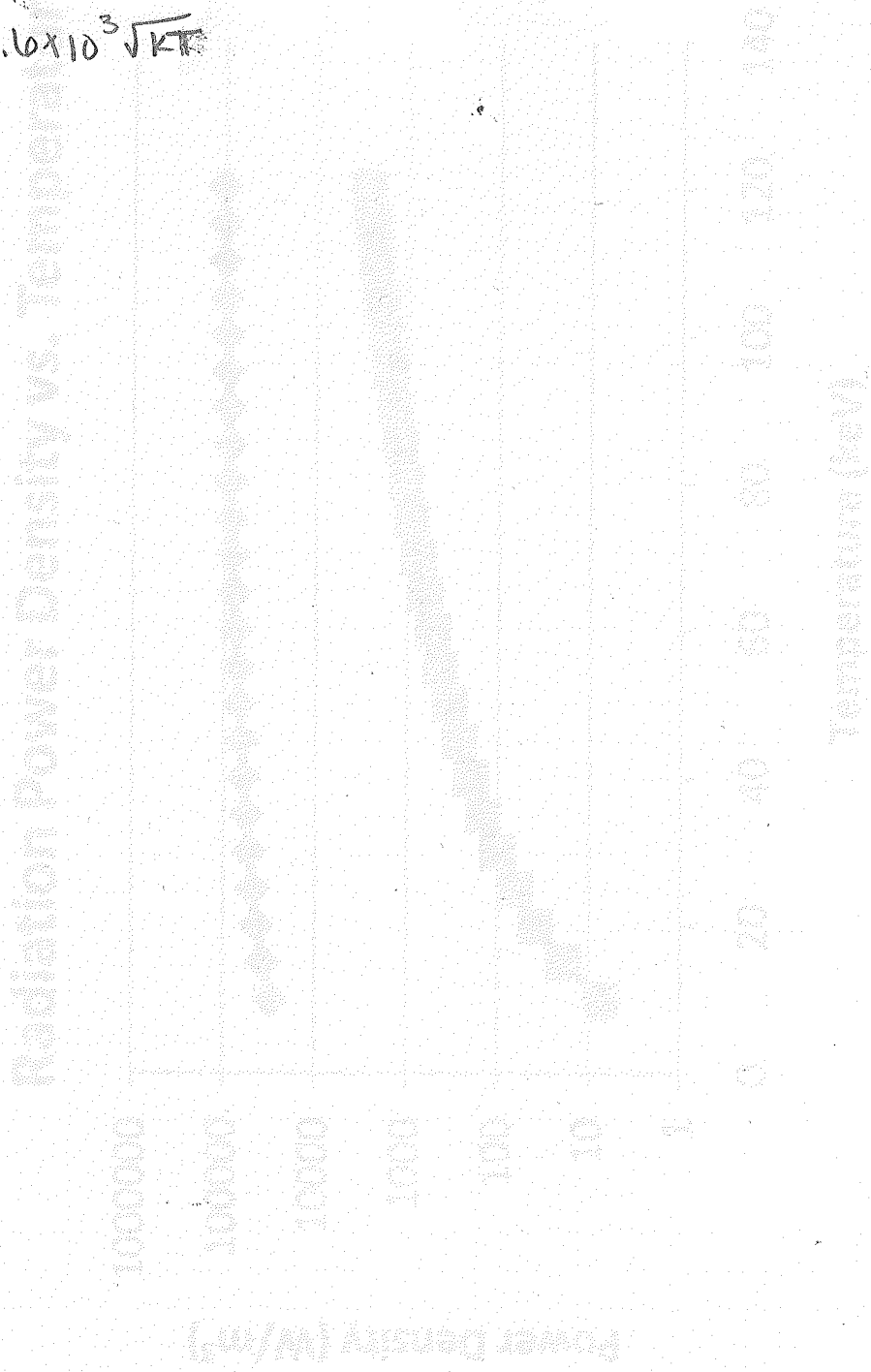
$$c) \Psi(T) \approx 10^{-4.2} [T(\text{keV})/1\text{keV}]^{1.4} \quad T \rightarrow 10-120\text{keV}$$

$$P_{\text{cyc}} = A_{\text{cyc}} N_c B^2 k T_e \Psi$$

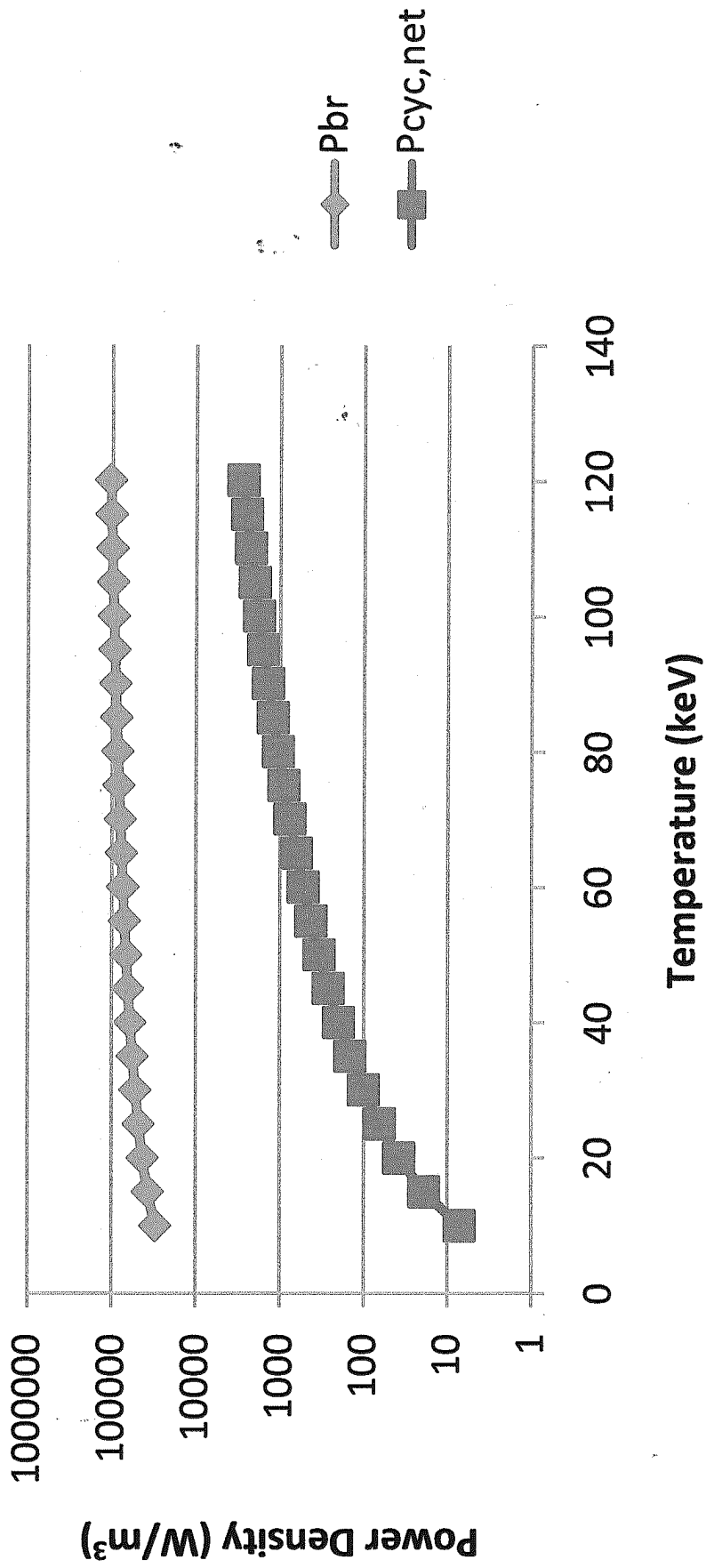
$$P_{\text{cyc}}^{\text{net}} = 453.6 k T \left[10^{-4.2} \left[\frac{k T}{1\text{keV}} \right]^{1.4} \right]$$

$$P_{\text{br}} = A_{\text{br}} N_d N_e Z^2 \sqrt{k T} + A_{\text{br}} N_h N_e Z^2 \sqrt{k T}$$

$$P_{\text{br}} = 9.6 \times 10^3 \sqrt{k T}$$



Radiation Power Density vs. Temperature



6.2 Transport equation (6.42)

$$\frac{\partial N}{\partial t} + \vec{v} \cdot \nabla_r N + \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v N = \left(\frac{\partial N}{\partial t}\right)_c + \left(\frac{\partial N}{\partial t}\right)_s$$

a) $q=0$, equation becomes

$$\frac{\partial N}{\partial t} + \vec{v} \cdot \nabla_r N = \left(\frac{\partial N}{\partial t}\right)_c + \left(\frac{\partial N}{\partial t}\right)_s$$

b) For monoenergetic particles, $E = \text{constant}$.

$$\Rightarrow v = \text{constant, i.e., } v_x^2 + v_y^2 + v_z^2 = \text{constant}$$

The transport equation becomes

$$\left\{ \begin{array}{l} \frac{\partial N}{\partial t} + \vec{v} \cdot \nabla_r N + \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v N = \left(\frac{\partial N}{\partial t}\right)_c + \left(\frac{\partial N}{\partial t}\right)_s \\ v_x^2 + v_y^2 + v_z^2 = \text{constant} = \frac{2E}{m} \end{array} \right.$$

c) only the z -direction is relevant, implying $\frac{\partial N}{\partial x} = \frac{\partial N}{\partial y} = 0$, $\frac{\partial N}{\partial v_x} = \frac{\partial N}{\partial v_y} = 0$

$$\vec{v} \cdot \nabla_r N = v_x \cdot \frac{\partial N}{\partial x} + v_y \cdot \frac{\partial N}{\partial y} + v_z \cdot \frac{\partial N}{\partial z} = v_z \cdot \frac{\partial N}{\partial z}$$

$$\vec{v} \times \vec{B} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ v_x & v_y & v_z \\ B_x & B_y & B_z \end{vmatrix} = (v_y B_z - B_y v_z) \hat{i} + (v_z B_x - v_x B_z) \hat{j} + (v_x B_y - v_y B_x) \hat{k}$$

$$\text{So, } (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v N = E_z \cdot \frac{\partial N}{\partial v_z} + (v_x B_y - v_y B_x) \cdot \frac{\partial N}{\partial v_z}$$

\Rightarrow Transport equation becomes

$$\frac{\partial N}{\partial t} + v_z \frac{\partial N}{\partial z} + \frac{q}{m} \left[E_z \frac{\partial N}{\partial v_z} + (v_x B_y - v_y B_x) \cdot \frac{\partial N}{\partial v_z} \right] = \left(\frac{\partial N}{\partial t}\right)_c + \left(\frac{\partial N}{\partial t}\right)_s$$

d) The equation is just the same as (6.42), since the source term and collision terms need to be considered

6.3. Plasma diffusion follows

$$\frac{\partial}{\partial t} N - D \nabla^2 N = 0$$

For the one-dimensional case, $\frac{\partial}{\partial t} N - D \frac{d^2 N}{dx^2} = 0$.

Let $N(t, x) = T(t) \cdot R(x)$. — separation of variables.

$$\frac{dT(t)}{dt} \cdot R(x) - D \cdot \frac{d^2 R(x)}{dx^2} \cdot T(t) = 0$$

$$\frac{1}{T} \frac{dT}{dt} = D \cdot \frac{1}{R} \frac{d^2 R}{dx^2}$$

Let each side equals to $-\lambda^2$ (it equals to a negative value since density decays w/ time)

$$\frac{1}{T} \frac{dT}{dt} = -\lambda^2 \Rightarrow T = C_1 e^{-\lambda^2 t}$$

$$D \cdot \frac{1}{R} \frac{d^2 R}{dx^2} = -\lambda^2 \Rightarrow R = C_2 \sin\left(\frac{\lambda}{\sqrt{D}} x + C_3\right)$$

Boundary conditions. $N(x=0) = N(x=x_0) = 0$.

$$\Rightarrow R(x=0) = R(x=x_0) = 0.$$

$$\Rightarrow \begin{cases} C_3 = 0 \\ \frac{\lambda}{\sqrt{D}} x_0 = \pi, 2\pi, 3\pi \end{cases}$$

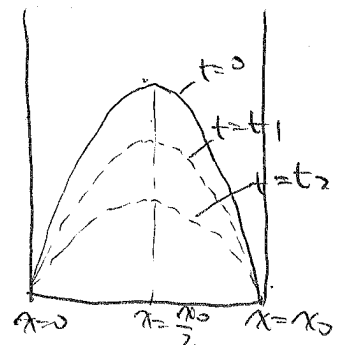
$$\hookrightarrow \lambda = \frac{\sqrt{D} \pi}{x_0}$$

$$N(t, x) = C \cdot \sin\left(\frac{\pi x}{x_0}\right) \cdot e^{-\frac{D \pi^2}{x_0^2} t}$$

Initial condition. $N(t=0, x=\frac{x_0}{2}) = N_0$.

$$\Rightarrow C = N_0$$

$$N(t, x) = N_0 \sin\left(\frac{\pi x}{x_0}\right) e^{-\frac{D \pi^2}{x_0^2} t}$$



6.4. With the source term, (6.13) becomes

$$\frac{\partial}{\partial t} N(\vec{r}, t) - D \nabla^2 N(\vec{r}, t) = S \quad \leftarrow \text{constant}$$

For steady state,

$$-D \nabla^2 N(\vec{r}, t) = S$$

$$-D \cdot \frac{1}{r} \frac{d}{dr} \left(r \frac{dN}{dr} \right) = S$$

$$\Rightarrow \frac{d}{dr} \left(r \frac{dN}{dr} \right) = -\frac{S}{D} r$$

$$\Rightarrow r \frac{dN}{dr} = -\frac{S}{2D} r^2 + C_1$$

$$\Rightarrow \frac{dN}{dr} = -\frac{S}{2D} r + \frac{C_1}{r}$$

$$\Rightarrow N = -\frac{S}{4D} r^2 + C_1 \ln r + C_2$$

BC(1) $N(r \rightarrow \infty) \neq \text{infinite} \Rightarrow C_1 = 0$.

BC(2) $N(r=0) = N(0) \Rightarrow C_2 = N(0)$

$$N = -\frac{S}{4D} r^2 + N(0)$$

\hookrightarrow We are not able to get (6.54) though.

6.5. You can either use the above relationship or (6.54).

Here I use (6.54).

$$\text{Average density} = \frac{\int N dV}{\int dV} = \frac{\int_0^a N \cdot r dr \int_0^{2\pi} d\theta \int_0^H dh}{\int_0^a r dr \int_0^{2\pi} d\theta \int_0^H dh}$$

$$= \frac{\int_0^a N(0) \sqrt{1 - \frac{r^2}{a^2}} \cdot r dr}{\frac{a^2}{2}}$$

$$= \frac{N(0) \left(1 - \frac{r^2}{a^2} \right)^{\frac{3}{2}} \cdot \left(-\frac{2}{3} a^2 \right) N(0) \Big|_0^a}{a^2/2}$$

$$= \frac{2}{3} N(0)$$

